



Exploiting Subradiance and Selective Radiance in Dense Atomic Arrays





M. Moreno-Cardoner UAB, ICFO (Barcelona, Spain)

D. E. Chang ICFO (Barcelona, Spain)

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Atom-Photon Interactions

A Single Atom

Ultimate non-linear optical element

non-classical states of light, quantum information processing



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Problem: A single atom and a single photon **do not like to talk** to each other

Atom-Photon Interactions

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Spontaneous Emission into undesired directions is a fundamental limitation: Source of **(quantum) information loss**

Error of most applications is bounded by the "Cooperativity" or "Optical Depth"

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 Γ_{1D}



$$OD = \frac{\Gamma_{1\mathrm{D}}}{\Gamma'} \times N$$

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Universal Approach to Optimal Photon Storage in Atomic Media

Single Photon Storage

Alexey V. Gorshkov,¹ Axel André,¹ Michael Fleischhauer,² Anders S. Sørensen,³ and Mikhail D. Lukin¹

off-resonant Raman fields to photon-echo-based techniques. Furthermore, we derive an optimal control strategy for storage and retrieval of a photon wave packet of any given shape. All these approaches, when optimized, yield identical maximum efficiencies, which only depend on the optical depth of the medium.

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Error ~ 1/OD

Quantum interface between light and atomic ensembles

Klemens Hammerer

Institute for Theoretical Physics, University of Innsbruck, and Institute for Quantum Optics and Quantum Information, Austrian Academy of Science, Technikerstrasse 25, 6020 Innsbruck, Austria

Anders S. Sørensen and Eugene S. Polzik

Niels Bohr Institute, Copenhagen University, Blegdamsvej 17, Copenhagen 2100, Denmark



The bound on the achievable squeezing is not so surprising, given that the state of atoms suffered essentially a decay by a fraction η_A . Due to the relation $\kappa^2 = d\eta_A$, cf. Eq. (59), there is always an optimal choice for the decay η_A given a certain optical depth *d* (Hammerer *et al.*, 2004). For the decoherence model adopted here the limit to spin squeezing by QND measurement and feedback is $\Delta P_{A,\text{final}}^2 \ge (1+\sqrt{1+d})^{-1} - d^{-1/2}$, for large optical density. For a detailed discussion on the limits of spin

Spin Squeezing

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Spin Squeezing

Error decreases as 1 / OD or slower















rich and engineerable emission pattern into 4 π





rich and engineerable emission pattern into 4 $\!\pi$







rich and engineerable emission pattern into 4 π



subradiance suppression of decay rate

 ϵ

 $\mathbf{r} \rightarrow \mathbf{E}(\mathbf{r}) = \mu_0 \omega^2 \mathbf{G}(\mathbf{r}, 0, \omega) \cdot \boldsymbol{\wp}$

field created by a single dipole is given by Green's function

scattered field





atomic dynamics effective Hamiltonian

$$\hat{H}_{\text{eff}} = -\mu_0 \omega_{eg}^2 \wp^2 \sum_{i,j} G(\mathbf{r}_i, \mathbf{r}_j, \omega_{eg}) \hat{\sigma}_i^{eg} \hat{\sigma}_j^{ge}$$



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Regular and **Dense Arrays** maximize the effects of Wave Interference





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a sub-wavelength chain can support single-excitation guided modes

> subradiant modes are guided

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a sub-wavelength chain can support single-excitation guided modes

subradiant modes are guided

A single layer of atoms ordered in a sub-wavelength array behaves as a perfect mirror



We can use such effects to create a *nearly perfect quantum memory*

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Exponential Improvement in Photon Storage Fidelities using Selective Radiance A. Asenjo-Garcia, M. Moreno-Cardoner, A. Albrecht, H. J. Kimble, D. E. Chang, *Phys. Rev. X* 7, 031024 (2017)

We can use such effects to create a *nearly perfect quantum memory*



"selectively radiant system"

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-xponential Improvement vs. Atom Number

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A small array of 4 x 4 atoms can be as efficient as a disordered ensemble with N = $10^6 \sim 10^7$



Optimization of photon storage fidelity in ordered atomic arrays M. T. Manzoni, M. Moreno-Cardoner, A. Asenjo-Garcia, J. V. Porto, A. V. Gorshkov, and D. E. Chang, New Journal of Physics, 20, 083048 (2018)

We can use such effects for an *efficient transport of excitations*



Extraordinary subradiance with lossless excitation transfer in dipole-coupled nano-rings of quantum emitters

M. Moreno-Cardoner, D. Plankensteiner, L. Ostermann, D. E. Chang, H. Ritsch, arXiv:1901.10598

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<u>Outlook</u>

 Including Wave Interference and Correlated Dissipation leads to new paradigms

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- Including Wave Interference and Correlated Dissipation leads to new paradigms
- Novel phenomena, More powerful protocols?



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Albrecht



T. Porto

(JQI)



A. Gorshkov (JQI)



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Α.

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Α.





A. Gorshkov (JQI)



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D. Plankensteiner (Univ. Innsbruck)



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Manzoni





(JQI)



A. Gorshkov (JQI)



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Α.
Albrecht





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